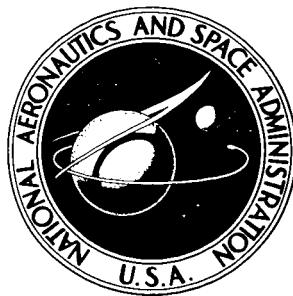


NASA TECHNICAL NOTE

NASA TN D-3913



N 67-23798

FACILITY FORM 602

_____ (ACCESSION NUMBER)
93 (PAGES)

_____ (THRU)
103 (CODE)
_____ (CATEGORY)

THE ANALYSIS AND CALIBRATION OF ANALOG TO DIGITAL ENCODERS

by J. A. Loger

Goddard Space Flight Center
Greenbelt, Md.

THE ANALYSIS AND CALIBRATION OF
ANALOG TO DIGITAL ENCODERS

By J. A. Loger

Goddard Space Flight Center
Greenbelt, Md.

NATIONAL AERONAUTICS AND SPACE ADMINISTRATION

For sale by the Clearinghouse for Federal Scientific and Technical Information
Springfield, Virginia 22151 - CFSTI price \$3.00

ABSTRACT

The conversion of continuously varying data into a form acceptable to a digital telemetry system introduces error into the data. This paper examines the errors in this signal conditioning process and derives expressions for them by using statistical methods. The standard calibration procedure for analog to digital encoders is examined and found to be inadequate. A method of calibration is presented which overcomes the inadequacies of existing procedures.

CONTENTS

Abstract	ii
INTRODUCTION	1
ANALYSIS OF THE ENCODER	2
THE CALIBRATION PROCEDURE	10
ERRORS IN THE TRANSDUCER-ENCODER SUBSYSTEM	16
CONCLUSIONS	18
ACKNOWLEDGMENT	18
References	19

THE ANALYSIS AND CALIBRATION OF ANALOG TO DIGITAL ENCODERS*

by

J. A. Loger

Goddard Space Flight Center

INTRODUCTION

The primary data channel of a satellite is usually some form of digital telemeter. The information presented to this channel generally originates either from discrete state input devices such as counters, relays, or digital logic, or from continuous transducers such as voltage dividers, current probes and temperature sensors. In the former, signal conditioning is an error-free process of matching the discrete states of the input to the discrete states of the data channel. In the latter case, signal conditioning introduces error into the data since it must convert the continuously varying data into discrete states (quantizing). Since there are usually several transducers, the signal conditioner must also sample and multiplex inputs from each into the one data channel.

This paper is concerned with the errors generated by the quantizing process. It attempts an analysis of the error contained in the theoretical model of the quantizer, and of the determination of error through the calibration of the actual device. Some space will also be devoted to the error introduced by reconversion of digitized data, at the receiving end of the channel, into their original form. Errors due to sampling and multiplexing of data are not considered in this paper since most of the quantities being sampled change very slowly compared to the sampling interval; either their harmonic content is not of much interest, or the variation can be determined by simple interpolation between data samples. For example, spacecraft battery temperature varies in terms of degrees per hour; with a sampling rate of 25 cps, the data appear virtually constant between samples in time.

Although the statistical equivalent of quantizing (rounding or grouping) was first examined by Sheppard (Reference 1) in 1898, the theoretical analysis of quantizers (analog to digital encoders) is more recent in origin. In 1948, Bennett (Reference 2) analyzed the power spectrum of the quantizing error; in 1956, Widrow (Reference 3) used mathematical statistics to examine an encoder model.

*Also submitted to the University of Maryland in partial fulfillment of the requirements for the degree of Master of Science.

The calibration of analog to digital encoders has not been covered in any detail in existing literature. Whatever methods are mentioned generally do not provide adequate measures of encoder errors, nor do they always lead to an accurate estimate of the encoder transfer function. This paper is a partial attempt to rectify this situation.

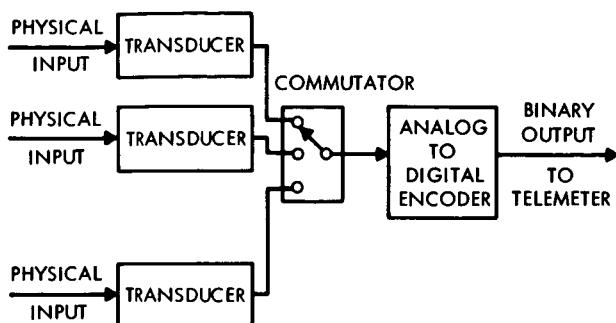


Figure 1—A transducer-encoder subsystem.

Transducer calibration usually consists of finding a least squares fit for a set of measured input-output data point pairs. This procedure is well covered in the literature of statistics (References 4, 5, and 6). The commutator is basically a multiple position electronic switch which selects different transducer outputs for input to the encoder. In this paper, the commutator is assumed to contribute no error to the signal.

ANALYSIS OF THE ENCODER

The encoder can be represented by a two-part network with a transfer characteristic which is a monotonic step function with a range of $2^n - 1$ steps, where n is the number of binary bits in the encoder output. The general relation between the input and output is

$$Q = kq, \quad \left(k - \frac{1}{2}\right) q < H(v) \leq \left(k + \frac{1}{2}\right) q, \quad k = 0, 1, 2, \dots, 2^n - 1, \quad (1)$$

where:

Q = binary output

q = quantizing interval

v = input in volts

$H(v)$ = a continuous function of the input.

The theoretical quantizer transfer function can be represented as:

$$Q = kq, \quad \left(k - \frac{1}{2}\right) q < x \leq \left(k + \frac{1}{2}\right) q, \quad k = 0, \pm 1, \pm 2, \dots, \quad (2)$$

where x is a continuous variable. Average quantizer gain is unity, whereas encoder gain is represented by a function of the input, $H(v)$, ideally of the form $H = mv$ where m is a constant. The encoder model can therefore be represented by a quantizer proceeded by an "amplifier" of gain $H(v)$ and $x = H(v)$. The term "amplifier" is used for convenience in the analysis even though the input voltage is converted to a continuous number rather than a voltage. Figure 2 is a block diagram of the analog to digital encoder model. The amplifier transfer function is shown in Figure 3a;

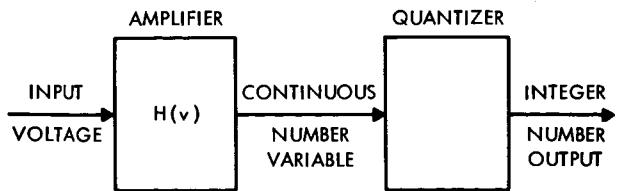
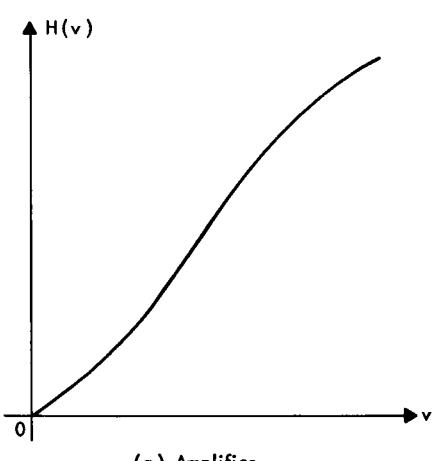
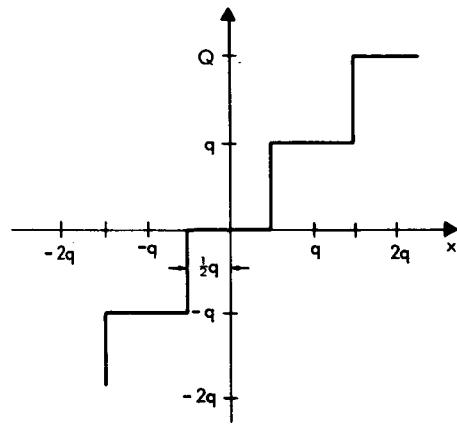


Figure 2—Model of analog to digital encoder.



(a) Amplifier



(b) Quantizer

Figure 3—Transfer functions.

the standard quantizer transfer characteristic is shown in Figure 3b. In essence, the quantizer produces an output which is the integer nearest to the value of the output. This process introduces a quantizing error into the data. This error is not of a random nature; for any given input, it is known exactly. In application, however, the quantizer input is not available to the system which receives the quantizer output and therefore the quantizing error cannot be evaluated for a given datum. Since the input is not available but is known to vary, the associated quantizing error over a large amount of data can be assumed to be random. This allows the use of statistical analysis as a means of evaluating the average effects of quantizer error on the quantizer output.

In a statistical analysis, the items of interest are usually the mean of a random variable given by

$$u_x = \int_{-\infty}^{\infty} x f(x) dx , \quad (3)$$

where

$$u_x = \text{mean of } x, \\ f(x) = \text{probability density function of } x(\text{p. d. f.}),$$

and the variance;

$$\sigma_x^2 = \int_{-\infty}^{\infty} (x - u_x)^2 f(x) dx . \quad (4)$$

The square root of the variance is the standard deviation σ_x . The mean u_x is the average value of the random variable x . In the case of error, it is desirable that the average value be zero. The standard deviation σ_x is the rms deviation of the random variable from its mean. It is a measure of the scattering of the data. For example, if the mean of the error is zero, then σ_x is the rms error. In general,

$$M_n = \int_{-\infty}^{\infty} (x)^n f(x) dx , \quad (5)$$

where M_n is the n th moment of the random variable x . The characteristic function (c.f.) of a probability density function is defined as

$$\phi_x(t) = \int_{-\infty}^{\infty} e^{itx} f(x) dx , \quad (6)$$

it has the property that

$$i^n M_n = \left. \frac{d^n \phi_x(t)}{dt^n} \right|_{t=0} , \quad (7)$$

where $i = \sqrt{-1}$ and the n th moment of $f(x)$ exists. This property will be used later to find the variance of a quantized variable. Equation 7 gives the moments around zero. σ_x^2 is a moment around u_x . The two moments are related by

$$\sigma_x^2 = M_2 - u_x^2 . \quad (8)$$

Also

$$\phi_x(0) = 1 . \quad (9)$$

This can be seen from Equation 6 and

$$\int_{-\infty}^{\infty} f(x) dx = 1. \quad (10)$$

The c.f. also has the property that if x and y are independent random variables, then:

$$\phi_{(x+y)}(t) = \phi_x(t)\phi_y(t). \quad (11)$$

If the input to the quantizer is a random variable, then the error produced by the quantizer must in some way alter its p.d.f. Examination of the input and output p.d.f.'s of the quantizer should reveal the effect.

Since any value of x in the range $kq \pm q/2$ is read as kq , the probability that the output of the quantizer is kq must be:

$$g(kq) = \int_{kq-q/2}^{kq+q/2} f(x) dx, \quad (12)$$

where $f(x)$ is the p.d.f. of the random variable x . The quantizer converts a continuous p.d.f. into a series of discrete probability pulses as shown in Figure 4. The function $g(kq)$ can be viewed as

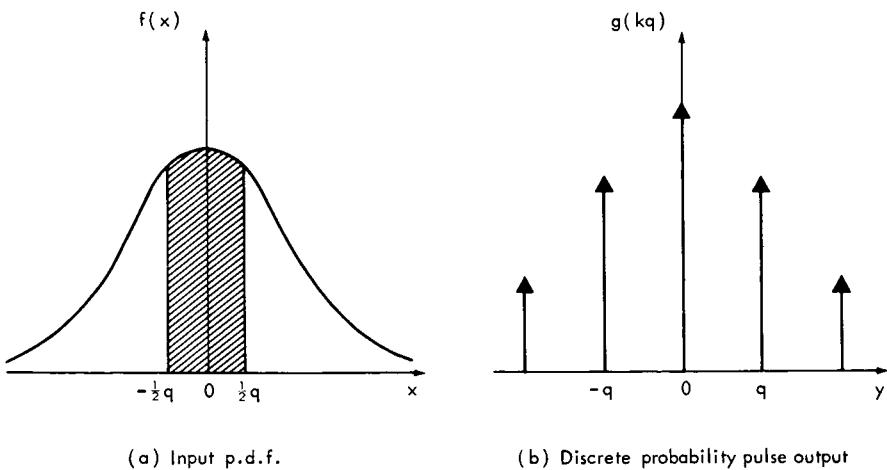


Figure 4—Functions relating to the quantizer.

a function $g(y)$ which has been sampled at $y = kq$, $k = \pm 1, \pm 2, \dots$. Now

$$\begin{aligned} g(y) &= \int_{y-q/2}^{y+q/2} f(x) dx \\ &= \int_{-\infty}^{\infty} f(x) U(y-x) dx, \end{aligned} \quad (13)$$

where

$$\begin{aligned} U(y-x) &= 1 \text{ for } y-q/2 < x \leq y+q/2 \\ &= 0 \text{ otherwise; } \end{aligned}$$

therefore

$$g^*(y) = \sum_k g(y) \delta(y - kq)$$

represents the sampled function and $\delta(y)$ is the Dirac-delta function. The sampled function $g^*(y)$ is the p.d.f. of the output if $f(x)$ is the p.d.f. of the input.

In its present form, $g^*(y)$ is too awkward to provide much information; however, the characteristic function of $g^*(y)$ will provide more information. This is true partly because of Equation 7, which states that if the moments of the p.d.f. of a random variable exist, they may be found by the various derivatives of the c.f. at the origin. Equation 13 is a case of the convolution integral

$$h(y) = \int_{-\infty}^{\infty} r(x) s(y-x) dx = \int_{-\infty}^{\infty} s(x) r(y-x) dx. \quad (14)$$

The important property of this integral is that the Fourier transform $\phi_h(t)$ of $h(y)$ is

$$\phi_h(t) = \phi_r(t) \phi_s(t). \quad (15)$$

The c.f. of $g^*(y)$ is

$$\begin{aligned}
 \phi_{g^*}(t) &= \int_{-\infty}^{\infty} e^{ity} g^*(y) dy \\
 &= \int_{-\infty}^{\infty} e^{ity} \left[\sum_{k=-\infty}^{\infty} g(y) \delta(y - kq) \right] dy \\
 &= \sum_{k=-\infty}^{\infty} e^{itkq} g(kq) .
 \end{aligned} \tag{16}$$

This relation is a Fourier series representation whose fundamental period is $2\pi/q$. Evidently, the characteristic function of the quantizer output is periodic with period $2\pi/q$. The waveform of $\phi_{g^*}(t)$ can now be found with the aid of Equation 15.

It can be shown (Reference 9) that the coefficients $g(kq)$ of the Fourier series are equal to the values of the Fourier integral of the typical waveform of $\phi_{g^*}(t)$ evaluated at kq , $k = \pm 1, \pm 2, \dots$, times the fundamental frequency. Thus

$$\phi_{g^*}(t) = \frac{1}{q} \sum_{k=-\infty}^{\infty} \phi_f \left(t - \frac{2\pi k}{q} \right) ;$$

but

$$\frac{1}{q} [g(y)] = \frac{1}{q} \int_{-\infty}^{\infty} f(x) U(y - x) dx ,$$

and therefore, by Equations 6 and 15,

$$\phi_{g^*}(t) = \sum_{k=-\infty}^{\infty} \phi_f \left(t - \frac{2\pi k}{q} \right) \frac{\sin \left(t - \frac{2\pi k}{q} \right) \frac{q}{2}}{\left(t - \frac{2\pi k}{q} \right) \frac{q}{2}} .$$

This relation does not appear to be better for analysis purposes than that shown in Equation 16. However, if a restriction is placed upon $\phi_f(t)$, it does lead to useful results. Equation 7 shows that the p.d.f. moments are determined by differentiating the c.f. at the origin. If the quantizing width q is made fine enough so that the c.f. of the input (the typical waveform) is zero for $|t| \geq 2\pi/q$, then there would be no contribution to the c.f. at the origin from other parts of the c.f. Consequently,

the derivatives of the c.f. at the origin would be the derivatives of

$$\phi_g(t) = \phi_f(t) \frac{\sin \frac{tq}{2}}{\frac{tq}{2}} . \quad (19)$$

By Equation 11, this is the characteristic function of the sum of two independent random variables, one of which is the input and the other is a uniform distribution of width q and uniform height $1/q$. This can be condensed into the following statement: If the characteristic function of a random input is zero for $|t| \geq 2\pi/q$, then the error due to quantization can be considered as an independent random variable with the uniform distribution described above, and the output is the sum of the two independent random variables, the quantizer input x , and quantization noise N .

The restriction implies that the p.d.f. of the input tapers off smoothly to zero at the ends of its range, since the Fourier transform of such a function has a definite "bandwidth." Taking the derivative of Equation 19 and letting

$$\phi_N(t) = \frac{\sin \left(\frac{tq}{2} \right)}{\frac{tq}{2}} .$$

we have

$$\frac{d\phi_{g*}(t)}{dt} = \frac{d\phi_x(t)}{dt} [\phi_N(t)] + \phi_x(t) \left[\frac{d\phi_N(t)}{dt} \right] .$$

By Equation 7, setting $t = 0$,

$$u_Q = u_x , \quad (20)$$

since $\phi_N(t) = 1$ at $t = 0$ and $d\phi_N(t)/dt = 0$ at $t = 0$. Since x and N are independent,

$$\sigma_Q^2 = \sigma_x^2 + \sigma_N^2 .$$

Using Equation 4,

$$\sigma_N^2 = \int_{-\infty}^{\infty} N^2 f(N) dN = \int_{-q/2}^{q/2} \frac{1}{q} N^2 dN = \frac{N^3}{3q} \Big|_{-q/2}^{q/2} = \frac{q^2}{12} . \quad (21)$$

Thus

$$\sigma_Q^2 = \sigma_x^2 + \frac{q^2}{12} \text{ and } \sigma_x^2 = \sigma_Q^2 - \frac{q^2}{12} . \quad (22)$$

Equations 20 and 22 are the first two of a series of correction formulas first derived by Sheppard (Reference 1) to find the various moments of the input distribution from the moments of the output of a quantizer. Widrow (Reference 3) analyzed the behavior of a quantizer with a Gaussian input. The results showed that for quantization intervals as large as twice the standard deviation of the input, the error in the calculated variance of the input, using Equation 22 is 3 percent. This error increased drastically for higher moments. The probability that a Gaussian variable deviates from its mean by more than 3σ is about .0027; 99.7 percent of the fluctuations are less than 3σ in magnitude; and 99 percent of the fluctuations are less than 2.6σ . This implies that, for an approximately Gaussian variable, useful estimates of mean square error can be made even if the range of the variations is only about 2.5 quantization levels. The errors in the corrections for Gaussian input are due to the fact that the c.f. of a Gaussian variable decreases in size as $e^{-t^2\sigma^2}$ and therefore never attains zero. It is used as an example to provide a feel for the size of quantization interval required for useful results.

In general, the narrower the range of the random variable, the wider the characteristic function. Consequently, a rule of thumb for the size of the quantization interval could be based upon the range of the random variable. If the random fluctuations in the input are roughly Gaussian, then a quantization width of one-third of the range is adequate. This leads to an error of only 3 percent in the estimation of the variance of the input. If the distribution of the input is known, a quantization width can be determined more accurately by examination of its characteristic function. This process is an attempt to insure the independence of the quantization noise and so make the error analysis of an encoder relatively simple. If such independence cannot be assumed, the error analysis may require more sophisticated and time-consuming techniques.

If the input to the quantizer consists of a sum of independent random variables of which only one satisfies the above criterion, then the sum satisfies the criterion. If the input $x = y + z + a$, where y , z and a are independent random variables, then by Equation 11:

$$\phi_x(t) = [\phi_y(t)][\phi_z(t)][\phi_a(t)] .$$

If y satisfied the criterion, then $\phi_y(t) = 0$ for $|t| \geq 2\pi/q$. Consequently, $\phi_x(t) = 0$ for $|t| \geq 2\pi/q$. Because of this property, it may be useful purposely to add noise of *known* amplitude to the analog data to aid in recovering the moments of the data. If, for example, the range of the input random variable were only one quantization interval, the input would not satisfy the conditions for the independence of quantization noise. If noise of rms amplitude $q/2$ were added to the input, the condition would be met and the variance of the input could be calculated. This technique will be illustrated in the section entitled "The Calibration Procedure."

Finally, if the distribution of the input to a quantizer is known, uniform quantization may not necessarily be the optimum form. One investigation of this topic (Reference 7) has shown that, for minimum quantization noise, some form of nonuniform quantization may be required. However, if the distribution is not known, the encoder system could consist of some form of adaptive quantizer which examines the input data and adjusts the quantization intervals accordingly.

THE CALIBRATION PROCEDURE

The results of the previous section were based on certain assumptions about the statistical properties of the quantizer. These properties are in turn based upon the average effects of a quantizer upon a large number of input data points. The theoretical quantizer is not a truly random device. For example, if the input to the quantizer is 4.3725 . . . then the output is 4.0000 . . . , always. This is in contrast to a truly random device such as an information channel whose output, for a given input symbol, has a finite probability of being any one of several symbols. During actual use, when only the output of the encoder is available, only statistical techniques are available for estimating the input and expressing error in the data. During calibration, however, this is not the case.

Calibration is a process for determining the real gain function of a device as opposed to the theoretical. In the previous section, the theoretical encoder was represented as a quantizer preceded by a linear amplifier. The real encoder can be represented by the same quantizer preceded by a nonlinear amplifier and perhaps a d.c. offset. The offset may be required to account for encoders which produce an output other than the midpoint of the quantization interval. The model for the actual encoder is shown in Figure 5.

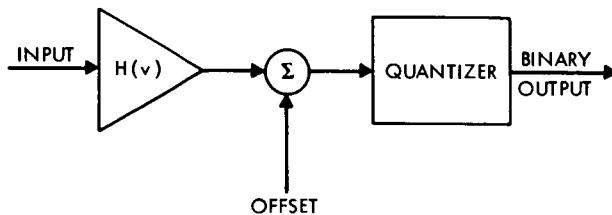


Figure 5—Encoder model.

Since the quantizer is assumed to be ideal, its transfer characteristic is known. Consequently, the purpose of the calibration is to determine the combined transfer characteristic of the amplifier and offset. This is complicated by the fact that the output of the latter two devices is observable only through the masking effect of the quantizer. In addition to the quantization error, there are fluctuations present in the real encoder transfer function due to variations in circuit parameters, changes in environment, etc.

In the model shown in Figure 5, these can be viewed as variations in $H(v)$ and the offset. As a result, the exact transfer function for the real encoder does not exist. There is, however, an average function around which the real one fluctuates. The purpose of calibration is to estimate this average curve. Calibration accuracy is measured by the closeness of the estimated average to the true one. The theory of estimation of functional relationships has been extensively treated in the literature of statistics (References 4, 5, and 6).

A method in general use for estimating the transfer function is the method of least squares. This technique minimizes the sum of the squares of the differences between the observed data and the estimated transfer function. In the case of the encoder, if $Y = H(x)$ is the transfer function of the amplifier, then the least squares estimate $H^*(x)$ of $H(x)$ is the one that minimizes

$$\sum_{j=1}^m [y_j - H^*(x_j)]^2 .$$

where (y_j, x_j) , $j = 1, 2, \dots, m$, are the observed output-input pairs. The method is based upon the Gauss-Markov theorem:

If the model for the device is expressible in the form $Y = X\theta + E$, where

$$Y = \begin{bmatrix} y_1 \\ y_2 \\ \vdots \\ y_m \end{bmatrix} \quad X = \begin{bmatrix} x_{11} & x_{12} & \dots & x_{1P} \\ x_{21} & & & \\ \vdots & & & \\ x_{m1} & & & x_{mP} \end{bmatrix} \quad \theta = \begin{bmatrix} \theta_1 \\ \theta_2 \\ \vdots \\ \theta_P \end{bmatrix} \quad E = \begin{bmatrix} e_1 \\ e_2 \\ \vdots \\ e_m \end{bmatrix},$$

with the average error $u_e = 0$, and if the errors e_j are independent of the y_j and have a common but unknown variance σ_e^2 , then the vector of coefficients $\hat{\theta}$ which minimizes $(Y - X\hat{\theta})' (Y - X\hat{\theta})$, where ' indicates transpose, is the linear, minimum variance unbiased estimator of θ . The term "unbiased" implies that $u_{\hat{\theta}} = \theta$; "minimum variance" means that the variance of $\hat{\theta}$ is less than or equal to the variance of any other estimator of θ , and the term "linear" means that the estimator of θ is expressible as a linear combination of the observed random variables Y .

This theorem allows the use of variety of mathematical models to represent the transfer characteristics of an electronic device. The only restrictions are (1) that the model be linear in the coefficients, and (2) that the errors in the observed outputs be statistically independent. The first condition permits models of polynomial form

$$Y = \sum_{k=1}^P \theta_k x^k,$$

and transcendental models of the form $Y = \theta_1 + \theta_2 e^{\pm x}$ and $Y = \theta \sin x$, but not models of the form $Y = \theta_1 + \theta_2 e^{\pm \theta_3 x}$. However, even the latter model can be transformed into an acceptable form. Restriction (2) above is satisfied in most calibration problems if some care is exercised. Since an encoder is a complex device, a fluctuation in the output is the sum of many variations within the equipment. As long as one of these variations is not unduly large, the errors in a given set of outputs can be assumed to be independent. Variables of the encoder which do cause large errors in the output must be controlled. For example, temperature and burn-in time have a considerable effect upon the encoder. For proper calibration, measurements must be made at different temperatures and times and transfer functions must be determined for each set of conditions.

The major drawback of the least squares method is that although given the model, it will determine the best estimates of the coefficients, it will *not* determine whether the model chosen is the best one for the data. Unfortunately, present methods of determining the best model are mostly intuitive. If polynomials are chosen, then the residual variance defined by

$$s_e^2 = \frac{1}{m-P} \sum_{j=1}^m \left[Y_j - \sum_{k=1}^P \theta_k x_j^{k-1} \right]^2$$

can be examined. A sufficiently large (~ 10) decrease in the residual variance due to using the next higher order polynomial can be taken as an indication that the higher order polynomial fits the data better. This is discussed in Reference 4. Even this technique can be misleading; for example, an $(m - 1)$ order polynomial will fit m data points exactly but, since the data points were observed with error, there is no assurance that the resulting polynomial is a good approximation of the underlying model.

The deficiencies in present encoder usage and calibration techniques stem either from misapplication of the least squares method, or from misunderstanding of the statistical properties of the encoder. The usual calibration process is fairly simple. An accurately known input voltage is stepped in equal increments through the range of interest, and for each input voltage an output binary number is recorded. The process is repeated several times; the outputs at each voltage are averaged over the several sweeps; and either a least squares fit is made of the inputs vs. averaged output pairs, or the averaged outputs are connected by straight line segments representing the transfer function interpolated between measured values. If there is a fluctuation in the binary output for any given input, the binary number which occurs most frequently is recorded as the output for that particular sweep. The calibration procedure is repeated at various temperatures (for example, -40°C , 25°C , 60°C), and the group of calibration curves is usually decorated with the cryptic remark " ± 1 count" or " $\pm \frac{1}{0}$ counts."

There are several reasons why this procedure is incorrect. First, the use of fixed input voltages may lead to a systematic error in calibration. In the model shown in Figure 5, the output corresponds to the midpoint of the rounding interval. There is, thus, a unique value of input which actually produces an amplifier output equal to this midpoint, but with no rounding error. The inputs in the calibration procedure generally do not correspond to the midpoints of rounding intervals. Consequently, the corresponding amplifier outputs are rounded off. The fluctuations in the amplifier and offset will cause a variation in the amount of roundoff, but the average round-off will not be zero. For example, if, for a given input, say 3.6500, the amplifier output is $72.647 \dots$, the output for the encoder is 73 and the rounding error is $.353 \dots$; fluctuations in the encoder may cause the amplifier output to vary, but the average of the variations will still have a rounding error of $.353 \dots$. The least squares method assumes that the average error in an output due to any input is zero. Since this assumption is not true, the calibration procedure introduces a bias into the measurement by associating the wrong input with the mid-point of the rounding interval. This illustrates the misapplication of the statistical properties of quantization noise.

Second, the use of fixed increment steps in the input violates the independence of errors condition of the Gauss-Markov theorem. The situation is illustrated in Figure 6; v_1 is applied to the input and Q_1 is observed as the encoder output. The amplifier output is Y_1 and the rounding error for the particular measurement is r_1 . The input is stepped by an amount Δv to v_2 , and the amplifier output is now Y_2 ; but $Y_2 = H(v_1 + \Delta v)$. If $H(v) = mv + b$ and $q = 1$, then $Y_2 - Y_1 = m\Delta v$, $r_2 = r_1 + (m\Delta v - 1)$ and, in general, $r_n = r_1 + (n - 1)(m\Delta v - 1)$. For the linear case, the rounding errors are a function of the output and therefore are not independent. This changes the values of the coefficients θ_k of the calibration curve. The same development could be used to show error dependence in higher order models.

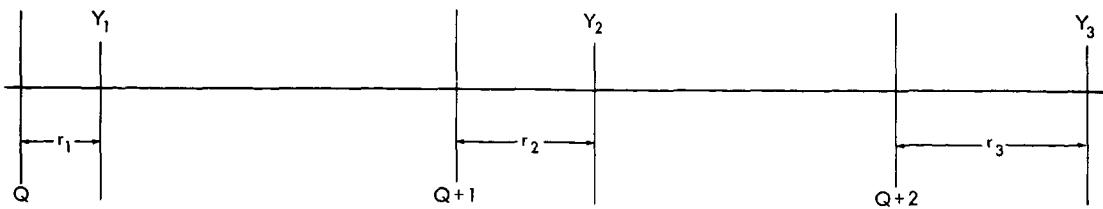


Figure 6—Use of fixed increment steps.

Finally, the choice of the most frequently occurring output truncates the range of the output random variable; the averages calculated from such data will therefore not be unbiased. This leads to even further error in the calculated calibration curve.

From the foregoing, it is apparent that the standard calibration scheme leads to input-output data pairs that do not provide a good estimate of the encoder transfer function. The calibration procedure is based on the assumption that, on the average, an output uniquely corresponds to an input. This is not so for the encoder. The corresponding assumption for use of the least squares method is that the underlying model is continuous and monotonic. The least squares method cannot be used for discrete and discontinuous models. The overall encoder transfer function clearly falls in this category.

It was mentioned previously that if the encoder is represented by a quantizer preceded by an amplifier, the only unknown in the model is the amplifier transfer function. This transfer function is continuous and monotonic. The offset is usually a known quantity since it is a function of encoder type and can be incorporated as a term in the amplifier transfer curve. The output of the encoder is the output of the amplifier rounded to the nearest integer. At a discrete set of inputs, the output of the encoder equals the output of the amplifier. Since the average effects of the fluctuations in the amplifier outputs are zero at these input points, the average error in the encoder output must be zero. If the input values were known, then the amplifier transfer curve would be known. In other words, at a given set of discrete input values, the transfer function of the encoder equals the transfer function of the amplifier. This obvious and very basic property of the encoder suggests a calibration method which avoids the pitfalls of the standard one.

For a given input v_j , the output of the amplifier is a quantity which is rounded to Q_j by the quantizer. Because of fluctuations within the amplifier, its output is occasionally rounded to some other, incorrect, number. The closer the amplifier's average output is to the edge of the rounding interval, the more frequently this occurs; it depends upon the probability that the size of the amplifier's output fluctuations exceeds the difference between the average output and the edge of the rounding interval.

If the fluctuations can be assumed to occur in both positive and negative directions with equal probability, the point of minimal variation in encoder output is when the average amplifier output is equal to the midpoint of the rounding interval. Also, at the midpoint, encoder output variations are positive and negative with about the same frequency. Thus, if the input to the encoder is varied

over an interval corresponding to a variation of about one rounding interval of the amplifier output, then the point of minimum and bidirectional variation in encoder output should be the midpoint of the rounding interval. The input and encoder output are now a data pair which can be used in determining the amplifier transfer curve. Strictly speaking, the procedure is estimating the median of the amplifier output fluctuations, not the midpoint of the rounding interval. The median ($u_{.5}$) is such that

$$\int_{-\infty}^{u_{.5}} f(x) dx = \int_{u_{.5}}^{\infty} f(x) dx$$

and is the point around which the fluctuations are positive and negative with equal probability. This corresponds to the midpoint of the rounding interval if the amplifier output p.d.f. is symmetrical. If the fluctuations can be assumed to be due to a large number of minute variations within the encoder, then this assumption is valid.

In actual measurement, because of the rounding effects of the quantizer, there will be some error in finding the coincidence of amplifier output and rounding interval midpoint. The input may correspond to an amplifier output which is not exactly equal to the midpoint. For small amplifier fluctuations there can be a small interval over which there is no appreciable fluctuation in the encoder output, and consequently the coincidence of amplifier output and rounding interval midpoint would be hard to identify. This could be avoided by arbitrarily adding a small amount of noise of zero average value to the input to increase the output fluctuations. In any case, (1) the error involved in the above procedure is much less than the size of the rounding interval; (2) on the average, the estimate of the midpoint should be correct; and (3) since the procedure is only concerned with estimating the midpoint of one rounding interval, the errors from one interval to the next are independent. Furthermore, by the formation of the model, the midpoints are parts of the continuous amplifier transfer curve. Any errors in determining the midpoint may be viewed as random errors around the midpoints. On the basis of the above, the data pairs satisfy the conditions for the application of the least squares method.

The question of exactly how the two procedures differ might arise since they both take associated input-output measurements of the encoder and fit them to a curve. Both use the discontinuous and discrete output of the encoder, yet one procedure permits the fitting of a continuous curve, and the other does not. The answer lies in the type of data being collected. The standard procedure does not recognize that the same output may be produced by a range of inputs, and consequently the data pairs obtained cannot determine a *unique* curve for the encoder. The other procedure attempts to find the input-output pairs for which the rounding error is zero; these are unique data pairs. Consequently these data pairs will determine a unique curve for the encoder.

One quantity associated with the encoder has not yet been discussed; the size of the fluctuations of the amplifier. In the normal calibration of a continuous device, this would have been

estimated by the square root of the residual variance mentioned previously. Unfortunately, the residual variance obtained by fitting a least squares curve to the data obtained by the above mentioned method, is not directly related to the size of the amplifier fluctuations. If the encoder input were set at one of the points determined to correspond to a rounding interval midpoint, the situation would correspond to the one described in the encoder analysis. The amplifier output would fluctuate with some distribution about the midpoint of the rounding interval. If the fluctuations were large enough, then the rounding errors could be assumed to be independent of the errors due to the fluctuations. And the variance of the amplifier output would be

$$\sigma_A^2 = \sigma_E^2 - \frac{q^2}{12} .$$

Since it is not known whether the amplifier output satisfies the independence criterion, noise, which is known to meet that criterion, must be added to the input. As was mentioned earlier, if any one of several independent random variables meets the criterion, then the sum of the random variables also meets the criterion. It will be shown later that the variance of the noise in the amplifier output due to the noise in the input is given by

$$\sigma_Y^2 = \left[\frac{dH(v)}{dv} \Big|_{v_j} \right]^2 \sigma_v^2 ,$$

where the derivative of the amplifier transfer curve is evaluated at the input value. Consequently, the variance of the output is

$$\sigma_E^2 = \sigma_Y^2 + \sigma_A^2 + \frac{q^2}{12} , \quad (23)$$

and therefore

$$\sigma_A^2 = \sigma_E^2 - \sigma_Y^2 - \frac{q^2}{12} . \quad (24)$$

The noise at the input can be obtained from a random noise generator and must have a zero average value. The rms value of the noise is σ_v . If the noise generator is roughly Gaussian, σ_v must be just large enough to produce a σ_Y of $q/2$, since a larger value might obscure the internal fluctuations of the amplifier. A number of samples are taken with the noise applied and the variance σ_E^2 is estimated by the formula

$$\hat{\sigma}_E^2 = S_E^2 = \sum_{j=1}^m \frac{(Q_j - \bar{Q})^2}{m-1}$$

where m is the number of samples, \bar{Q} is the average of the outputs, and S_E^2 is the estimator of the encoder variance. The number of samples required is not critical; about 30 should be sufficient. The number obtained by the use of Equation 24 should be a useful indicator of the mean square error in the encoder output due to the analog portions of the encoder. It would be useful to repeat this process at several points along the calibration curve and to take an average of the variance at these points.

The calibration curve should be accompanied by the estimated variance, not by the presently used statement of the range of variation. This allows the experimenter to judge the effects of the encoder upon his data. The statement ± 1 count does not provide information about the rms error in the encoder and, consequently, does not allow comparison of the encoder error with errors in other parts of the system.

During use, the amplifier calibration curve can be used as the transfer function for the encoder, and the effects of quantization can be considered as part of the total error of the subsystem.

ERRORS IN THE TRANSDUCER-ENCODER SUBSYSTEM

Several other sources of error exist in the transducer-encoder subsystem. Encoder error is only a part of the total error. For example, a transducer with an accuracy of ± 1 percent full-scale is a good transducer. Suppose that the ± 1 percent figure means that an error of 1 percent of the output represents a deviation of 3° from the average. Assuming that the encoder is perfect and only has a rounding error, the rms transducer error of 0.33 percent would produce an error component of 0.33 percent of full-scale of the encoder output. If the encoder is an 8-bit encoder, full-scale is 255, and the quantizing mean squared error is $[(100/255) \times \sqrt{1/12}]^2 = (.013\%)^2$ of full-scale. Since the transducer satisfies the quantizing criterion, the two errors are independent and the variances can be added. Therefore, the mean squared error of the subsystem output is $.111 + .013 = .124\%^2$. Obviously, the error due to the transducer is a significant part of the total error.

The above example uses percent error as a means to an end. The transfer of input error to the output of a device is not usually that simple. Only the assumption of a perfect encoder—i.e., amplifier function linear and invariant—allowed the use of percentages. In general, if the output of a device is related by a function, $v = F(x)$, to the input, then an error in an input x_1 of magnitude Δx produces a change in the output of $e = F(x_1 + \Delta x) - F(x_1)$. Now

$$\frac{e}{\Delta x} = \frac{F(x_1 + \Delta x) - F(x_1)}{\Delta x}$$

and if Δx is small,

$$\frac{F(x_1 + \Delta x) - F(x_1)}{\Delta x} \approx \left. \frac{dF(x)}{dx} \right|_{x_1}$$

thus

$$e \approx \left[\frac{dF(x)}{dx} \Big|_{x_1} \right] \Delta x \quad \text{and} \quad e^2 \approx \left[\frac{dF(x)}{dx} \Big|_{x_1} \right]^2 \Delta x^2 .$$

By the definition of variance,

$$\sigma_e^2 = \left[\frac{dF(x)}{dx} \Big|_{x_1} \right]^2 \sigma_{\Delta x}^2 \quad \text{and} \quad \sigma_e = \left[\frac{dF(x)}{dx} \Big|_{x_1} \right] \sigma_{\Delta x} . \quad (25)$$

The difference $F(x + \Delta x) - F(x)$ can be expanded as a Taylor series around x_1 . If Δx is small and the higher derivatives of $F(x)$ are small, then the remaining terms in the expansion are insignificant and the relations hold.

Equation 25 can be used to express the error present in the output of an encoder due to the effects of the transducer-encoder subsystem. For a given input x_j to the transducer with output v_j and a variance σ_T^2 , the expected mean squared error in the output of an encoder with analog transfer curve $H(v)$ and variance σ_A^2 is:

$$\sigma_E^2 = \left[\frac{dH(v)}{dv} \Big|_{v_j} \right]^2 \sigma_T^2 + \sigma_A^2 + \frac{q^2}{12} . \quad (26)$$

This expression shows that the error in the encoder output may partly depend on the magnitude of the input. If so, then errors of a certain percentage of full scale in the input will not produce the same percentage error in the output. For this reason, it is best to transpose all errors to the output before converting to percent full-scale. Equation 26 shows that the rounding error is only one part of the total error associated with the encoder output.

The information contained in the output of an encoder is not generally useful in its binary form. Consequently, it is converted to the physical units of the input transducer. This requires the inversion of the encoder and the transducer transfer functions and produces the relations $v_j = H^{-1}(Q_j)$ and $x_j = F^{-1}(v_j)$ where $v = F(x)$ is the transfer function of the transducer. By Equation 25, the variance of the encoder input v_j as calculated from the encoder output at Q_j is:

$$\begin{aligned} \sigma_{v_j}^2 &= \left[\frac{dH^{-1}(Q)}{dQ} \Big|_{Q_j} \right]^2 \left(\sigma_A^2 + \frac{q^2}{12} \right) \\ &= \frac{\left(\sigma_A^2 + \frac{q^2}{12} \right)}{\left[\frac{dH(v)}{dv} \Big|_{v_j} \right]^2} . \end{aligned} \quad (27)$$

Adding the transducer variance and again using Equation 25 gives the variance associated with the value of the input to the subsystem as calculated from the encoder output:

$$\sigma_x^2 = \frac{\sigma_{v_j}^2 + \sigma_t^2}{\left[\frac{dF(x)}{dx} \Big|_{x_j} \right]^2} \quad . \quad (28)$$

x_j and v_j are the estimated values of the input as determined from the calibration curves.

The transfer function of the encoder will usually be linear enough, for the purposes of calculating σ_x^2 , so that $dH(v)/dv$ can be assumed constant. This does not apply generally to the transducer transfer curve and, as a result, error estimates should be made at several points.

The quantities used in actual calculation are only estimates of the underlying model parameters and therefore the results are only approximate. However, they do provide an estimate of the size of error in both the observed output data and the calculated input of the transducer-encoder subsystem.

CONCLUSIONS

If certain conditions are satisfied, the errors produced by the rounding in an analog-to-digital encoder may be assumed to be independent of the input. This assumption leads to simple expressions for the total error contained in the output of a transducer-encoder subsystem. The errors contained in the input value, as calculated from the observed output value, depend upon the transfer characteristics of both the encoder and the transducer.

The encoder transfer curve cannot be determined accurately with standard calibration techniques. Such techniques lead to errors in the calibration data which violate the assumptions required for regular curve fitting methods. In this paper a calibration procedure has been derived which overcomes these difficulties and permits the use of regular curve fitting methods.

ACKNOWLEDGMENT

The author wishes to express his appreciation to Dr. R. M. Ginnings of the Department of Electrical Engineering, University of Maryland, for his invaluable advice and encouragement.

Goddard Space Flight Center
National Aeronautics and Space Administration
Greenbelt, Maryland, October 16, 1966
841-11-78-01-51

REFERENCES

1. Sheppard, W. F., "On the Calculation of the Most Probable Values of Frequency Constants for Data Arranged According to Equidistant Divisions of a Scale," *Proc. London Math. Soc.*, 29:353-380, 1898.
2. Bennett, W. R., "Spectra of Quantized Signals," *Bell System Tech. J.* 27(3):446-476, July, 1948.
3. Widrow, B., "A Study of Rough Amplitude Quantization By Means of Nyquist Sampling Theory," *I.R.E. Trans. on Circuit Theory*, PGCT-3(4):266-276, December, 1956.
4. Graybill, F. A., "An Introduction to Linear Statistical Models," Vol. 1, New York: McGraw-Hill, 1961.
5. Mandel, J., "The Statistical Analysis of Experimental Data," New York: John Wiley and Sons, 1964.
6. Natrella, M. G., "Experimental Statistics," National Bureau of Standards Handbook 91, 1963.
7. Morakis, J. C., "Optimum Quantizations," Goddard Space Flight Center Document X-520-65-274, July, 1965.
8. Kendall, M. G., Stuart, A., "The Advanced Theory of Statistics," Vols. I and II, New York: Hafner, 1948.
9. Papoulis, A., "The Fourier Integral and Its Applications," New York: McGraw-Hill, 1962.
10. Susskind, A. K., "Notes on Analog-Digital Conversion Techniques," New York: John Wiley and Sons, 1957.

10-3-67

"The aeronautical and space activities of the United States shall be conducted so as to contribute . . . to the expansion of human knowledge of phenomena in the atmosphere and space. The Administration shall provide for the widest practicable and appropriate dissemination of information concerning its activities and the results thereof."

—NATIONAL AERONAUTICS AND SPACE ACT OF 1958

NASA SCIENTIFIC AND TECHNICAL PUBLICATIONS

TECHNICAL REPORTS: Scientific and technical information considered important, complete, and a lasting contribution to existing knowledge.

TECHNICAL NOTES: Information less broad in scope but nevertheless of importance as a contribution to existing knowledge.

TECHNICAL MEMORANDUMS: Information receiving limited distribution because of preliminary data, security classification, or other reasons.

CONTRACTOR REPORTS: Scientific and technical information generated under a NASA contract or grant and considered an important contribution to existing knowledge.

TECHNICAL TRANSLATIONS: Information published in a foreign language considered to merit NASA distribution in English.

SPECIAL PUBLICATIONS: Information derived from or of value to NASA activities. Publications include conference proceedings, monographs, data compilations, handbooks, sourcebooks, and special bibliographies.

TECHNOLOGY UTILIZATION PUBLICATIONS: Information on technology used by NASA that may be of particular interest in commercial and other non-aerospace applications. Publications include Tech Briefs, Technology Utilization Reports and Notes, and Technology Surveys.

Details on the availability of these publications may be obtained from:

**SCIENTIFIC AND TECHNICAL INFORMATION DIVISION
NATIONAL AERONAUTICS AND SPACE ADMINISTRATION
Washington, D.C. 20546**